



## 2D semiconductor SnP2S6 as a new dielectric material for 2D electronics

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## ARTICLE

## 2D semiconductor $\text{SnP}_2\text{S}_6$ as a new dielectric material for 2D electronics

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Due to the intriguing optical and electronic properties, 2D materials are promising for next generation optoelectronic and electric device applications. Discovering new 2D materials with novel physical properties are rewarding for this area. In this work, we systematically investigated the optoelectronic properties of 2D metal thiophosphate  $\text{SnP}_2\text{S}_6$  with unique nanoporous structure. The intermediate bandgap makes  $\text{SnP}_2\text{S}_6$  a good candidate as both the channel and gate dielectric materials in the transistor device.  $\text{SnP}_2\text{S}_6$  showed good photoresponse properties. In addition, the  $\text{MoS}_2$  transistor with  $\text{SnP}_2\text{S}_6$  as dielectric layer shows a high dielectric constant ( $\approx 23$ ), low subthreshold slope down to 69.4 mV/dec, and it presented a 0.1 pA scale leakage current, a threshold voltage as low as 1.1 V, a ON/OFF ratio reaching  $10^7$  and negligible hysteresis with high stability and reproducibility. This work would open up new avenues for the discovery of new metal thiophosphate systems for future device applications.

### INTRODUCTION

The emergence of two-dimensional materials has brought new opportunities to the scientific and industrial fields. Due to the atomic ultrathin thickness, the quantum confinement effect in the thickness direction of two-dimensional materials is prominent, which demonstrated advantages of high carrier mobility, channel regulation capability and rich properties in force, heat, sound, light, electricity and magnetism<sup>1-4</sup>. As a result, novel electronic and optoelectronic devices based on 2D materials are developed and studied, and this is an active research field still today<sup>5-7</sup>. The rapid development of 2D materials is beneficial from the discovery of novel new 2D materials with intriguing physical properties. For example, the discovery of 2D  $\text{CuInP}_2\text{S}_6$  brings the 2D ferroelectricity and related device application into the center of ferroelectric transistor and memory device field. The isolation of layered  $\text{CrI}_3$  enables truly 2D ferromagnetic properties, attracting a lot of attention in the study of 2D magnetism. Layered ternary compounds such as  $\text{Ta}_2\text{PdS}_6$ <sup>8</sup>,  $\text{In}_2\text{P}_3\text{S}_9$ <sup>9</sup>,  $\text{Bi}_2\text{O}_2\text{Se}$ <sup>10, 11</sup>, also provide new opportunities for the electronic and optoelectronic performance improvement for 2D

material applications. There exists huge potential, opportunities as well as challenges for exploring new 2D material systems for the applications in future nanoelectronics and optoelectronics<sup>12, 13</sup>.

For dielectric application,  $\text{SiO}_2$  is used as dielectric layer typically, but traditional silicon-based field effect transistors (FETs) face various challenges when the device reduce to nanometer scale, the most prominent of which are reduced mobility and increased short-channel effects<sup>14</sup>. In order to mitigate these problems, novel dielectrics with remarkable properties begin to grow up. Advanced functional materials based on eco-friendly cellulose and sputtered multicomponent dielectrics can be processed under low temperature, and were reported to be implemented as high-capacitance gate dielectrics, which brings new insights into flexible and low-cost transistors and the need to meet the constraints for better band-offset matching<sup>15, 16</sup>. Moreover, 2D insulators are regarded as the ultimate solution for dielectric miniaturization<sup>17</sup>. Hexagonal boron nitride (h-BN) has been explored as the most promising 2D vdW dielectric layer, while novel 2D materials suitable for dielectric layer are rarely reported. However, the low dielectric constant ( $\approx 5.0$ ) and excessive leakage current make h-BN unsuitable for ultrascaled FETs with low power consumption<sup>18, 19</sup>. Therefore, it's important to find new materials to serve as reliable gate dielectric in FETs.

Among the various 2D material systems, the novel 2D metal thiophosphate  $\text{MP}_2\text{X}_6$  (M = metal, X = S or Se)<sup>20</sup> has received much attention due to their intriguing physical properties, including magnetic, electronic, ferroelectric, and optical characteristics<sup>21-23</sup>. Among the most extensively studied systems contain the  $\text{Sn}_2\text{P}_2\text{S}_6$  crystal with considerable piezo effect, proper ferroelectric and nonlinear optic properties<sup>24, 25</sup>,  $\text{MPS}_3$  as UV photodetector<sup>26</sup> which possess magnetic order<sup>27</sup>, and  $\text{CuInP}_2\text{S}_6$  heterostructure for non-

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volatile ferroelectric switches, memory devices and field-effect transistors<sup>28, 29</sup>.  $\text{SnP}_2\text{S}_6$ , a stable semiconductor material with a sizable indirect band gap of 2.23eV<sup>30</sup>, is also of importance to be investigated as a member of this family.  $\text{SnP}_2\text{S}_6$  exhibits a unique nonporous structure as it lacks half of the metal ions in the crystal structure of  $\text{SnP}_2\text{S}_6$  compared to the parent  $\text{Sn}_2\text{P}_2\text{S}_6$  structure. Due to the intrinsic 2D nanoporous structure, these days has witnessed more theoretical researches in  $\text{SnP}_2\text{S}_6$ <sup>24, 30-32</sup>, despite the recent experimental study of strong nonlinear optical response related with its inversion symmetry broken structure<sup>33, 34</sup>. Recent theoretical study has predicted that  $\text{SnP}_2\text{S}_6$  could have high carrier mobility<sup>20</sup>,

good switching stability, excellent insulation performance and gate-control ability. The measured dielectric constant is higher than the extensively used h-BN substrate<sup>35</sup>. Thus, it demonstrates great potential to be used as a channel or dielectric layer material in field effect transistors.

In this work,  $\text{SnP}_2\text{S}_6$  is firstly proposed as dielectric layer, the electronic and optoelectronic characteristics of which have been systematically investigated. Field effect transistors with excellent performance are successfully constructed, which verifies the feasibility of  $\text{SnP}_2\text{S}_6$  as dielectric layer material and provides new reliable building blocks for electronic devices based on 2D materials.

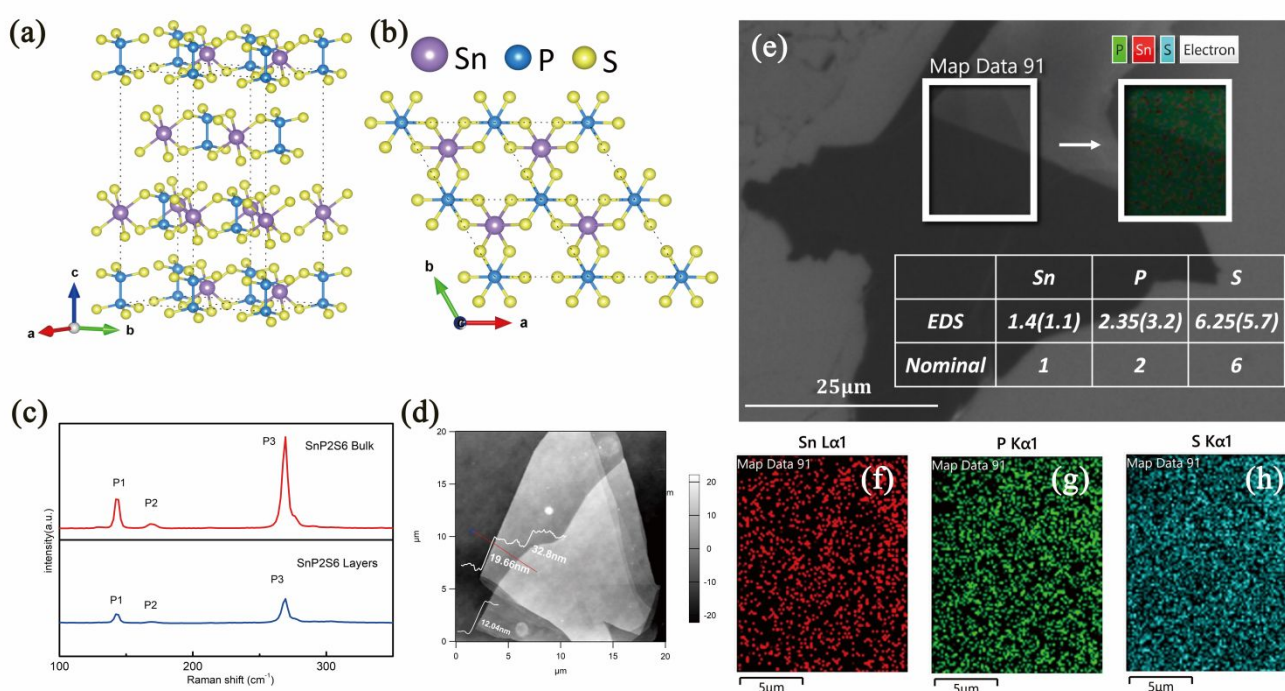


Fig 1 | Characterization of 2D material  $\text{SnP}_2\text{S}_6$ .

(a,b) Crystal structure of layered  $\text{SnP}_2\text{S}_6$ . (a) Side and (b) Top views. Purple, blue and yellow spheres represent Sn atoms, P atoms and S atoms, respectively. (c) Raman spectra of bulk and few-layer  $\text{SnP}_2\text{S}_6$ . (d) AFM topography of different-thickness  $\text{SnP}_2\text{S}_6$ . The thicknesses are around 12.04, 19.66 and 32.80 nm. (e) EDS results and elemental mappings of Sn (f), P (g) and S (h) for typical  $\text{SnP}_2\text{S}_6$  crystal, which shows homogenous distribution of the three elements.

## Materials characterization

Layered  $\text{SnP}_2\text{S}_6$  shows a rhombohedral layered structure (as shown in Fig.1(a, b)), which belongs to the R3 space group (No. 146) with rhombohedral symmetry<sup>30, 36</sup>, and only contains one formula unit<sup>20</sup> of one metal cation  $\text{Sn}^{4+}$  and one anionic  $[\text{P}_2\text{S}_6]^{4-}$ . As depicted in Fig.1(a, b), the ternary  $\text{SnP}_2\text{S}_6$  has a van der Waals layered structure and the interlayer spacing is about 0.65 nm, as reported in the literature<sup>33</sup>. Raman spectra were collected using a 532 nm laser in ambient conditions. As shown in Fig.1(c), three distinct peaks  $\text{P}_1$  ( $\approx 142 \text{ cm}^{-1}$ ),  $\text{P}_2$  ( $\approx 168 \text{ cm}^{-1}$ ) and  $\text{P}_3$  ( $\approx 269 \text{ cm}^{-1}$ ) can be identified in

both bulk and layered  $\text{SnP}_2\text{S}_6$ . Specifically, the three peaks are related to the internal stretching vibrations of S-P-S bonds<sup>9</sup>, which are consistent with the previous report<sup>33, 37</sup>. As the thickness increases, the intensity alters while positions of all peaks hardly change, which might be due to the weak interlayer interaction. In addition, as shown in Fig.1(e), it is noted that the atomic ratio of  $\text{SnP}_2\text{S}_6$  sample is 14: 23.5: 62.5 and 11: 32: 57 (totals 100), which is close to the stoichiometric ratio of 1:2:6 and the uniform distribution of elements Sn, P, and S are exhibited in Fig.1(f-h), as confirmed by Energy Dispersive X-ray Spectroscopy characterization (details in fig.S1).

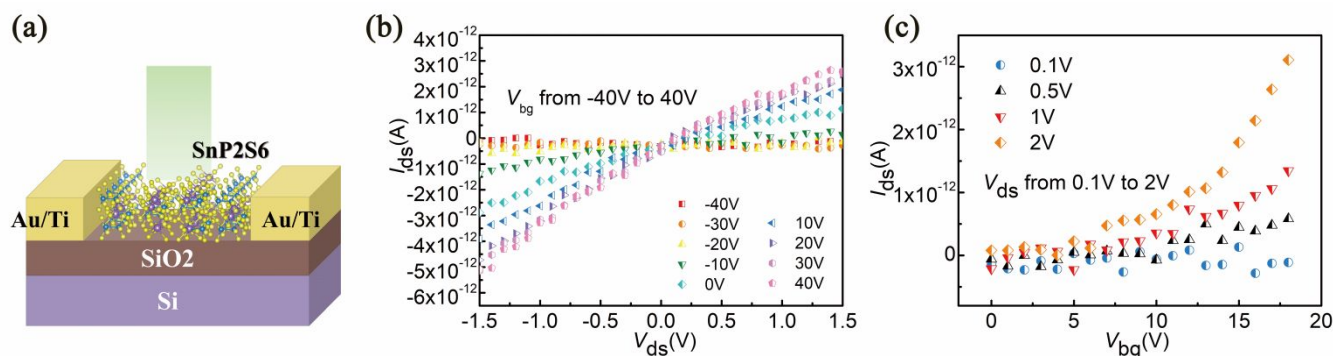


Fig 2 | Electronic properties of back-gated FETs using SnP<sub>2</sub>S<sub>6</sub> as the channel.

(a) Schematic of the SnP<sub>2</sub>S<sub>6</sub> photodetector on Si substrate, which provides the convenience of back gating. (b) Output curves ( $I_{ds}$ - $V_{ds}$ ) measured under  $V_{bg}$  from -40 V to 40 V. (c)  $I_{ds}$ - $V_{bg}$  curves of SnP<sub>2</sub>S<sub>6</sub> with different bias from 0.1 V to 2 V.

### Optoelectronic properties of SnP<sub>2</sub>S<sub>6</sub> thin flakes

SnP<sub>2</sub>S<sub>6</sub> phototransistor is fabricated on the standard SiO<sub>2</sub>/Si substrate (see methods part for details), and the schematic illustration of device is shown in Fig.2(a). The transfer and output curve of the transistor based on SnP<sub>2</sub>S<sub>6</sub> as the channel material are shown in Fig.2(b), (c). A linear relationship between  $I_{ds}$  and  $V_{ds}$  was observed. However, the transistor shows a very low electron mobility, the mobility ( $\mu = Ld/W\varepsilon_0\varepsilon_rV_{ds} \times dI_{ds}/dV_{bg}$ ) under a drain bias of 2 V is calculated as  $2.17 \times 10^{-5} \text{ cm}^2\text{V}^{-1}\text{s}^{-1}$ . The electronic performance of the device with various thicknesses of SnP<sub>2</sub>S<sub>6</sub> were also studied (detailed in Supplementary fig.S2 and fig.S3), demonstrating a similar insulating behaviour. The good insulating behaviour makes SnP<sub>2</sub>S<sub>6</sub> suitable for a dielectric layer, as discussed in details in next section.

The  $I_{ds}$ - $V_{ds}$  curves shown in Fig.3(a) are linear and symmetric for small bias voltages, indicating an ohmic like contact. The device shows an increase of drain current by several orders of magnitude as the device is illuminated. The typical optical image of the devices are shown in inset of Fig.3(a) and in fig.S2. Furthermore, the power dependence of the photocurrent is plotted with a log-log scale in Fig.3(b). As the incident light power increases from 0.1 to 50 mW, the photocurrent  $I_{ph}$  increases monotonically from  $8.3 \times 10^{-12}$  to  $1.2 \times 10^{-10}$  A with a bias of 5 V, indicating that the photocurrent is

linearly proportional to the incident light power following:  $I_{ph} \sim P^\beta$ . Moreover, as the gate voltage shifts from 0.1 to 5 V, the value of  $\beta$  increases from 0.23 to 0.44. The great loss of photocurrent represented by non-unity exponent relationship and low value of  $\beta$  suggest that the photocurrent is influenced by complex processes besides the absorption of photon, the generation of free carriers, recombination of the photogenerated free carriers, charge trapping by the defects and the charge impurities presented in SnP<sub>2</sub>S<sub>6</sub> flake, scarcely lateral photoeffect<sup>38</sup> and the adsorbed molecules at SnP<sub>2</sub>S<sub>6</sub>/SiO<sub>2</sub> interface due to the large surface-to-volume ratio<sup>39-41</sup>.

Meanwhile, the photoresponsivity  $R$  and photo-gain  $G$  at a drain voltage of  $V_{ds} = 5$  V exhibit strong dependence on the illumination power, as shown in Fig.3(c). The photoresponsivity  $R$  is one of the most important features for a photodetector<sup>42</sup>.  $R = I_{ph}/PS$ , where  $P$  is the incident light power density and  $S$  is the effective illuminated area. As shown in Fig.3(c), the photoresponsivity increases with the decrease of light power, which is possibly attributed to the less frequent carrier recombination and longer carrier lifetime under weaker illumination. This is consistent with photoconductor-based photodetector previously reported<sup>8, 43</sup>. Photogain ( $G$ ) is another parameter to evaluate the performance of a phototransistor, which can be calculated by the formula:  $G = (I_{ph}/q)/(PS/h\nu) = h\nu R/q$ ,<sup>44</sup> where  $I_{ph}$  is the photocurrent,  $q$  is the electron charge,  $h$  is Planck's constant,  $\nu$  is the light frequency, and  $R$  is responsivity. The value of  $G$  was estimated to reach up to 3.74 under the power of 0.1 mW.

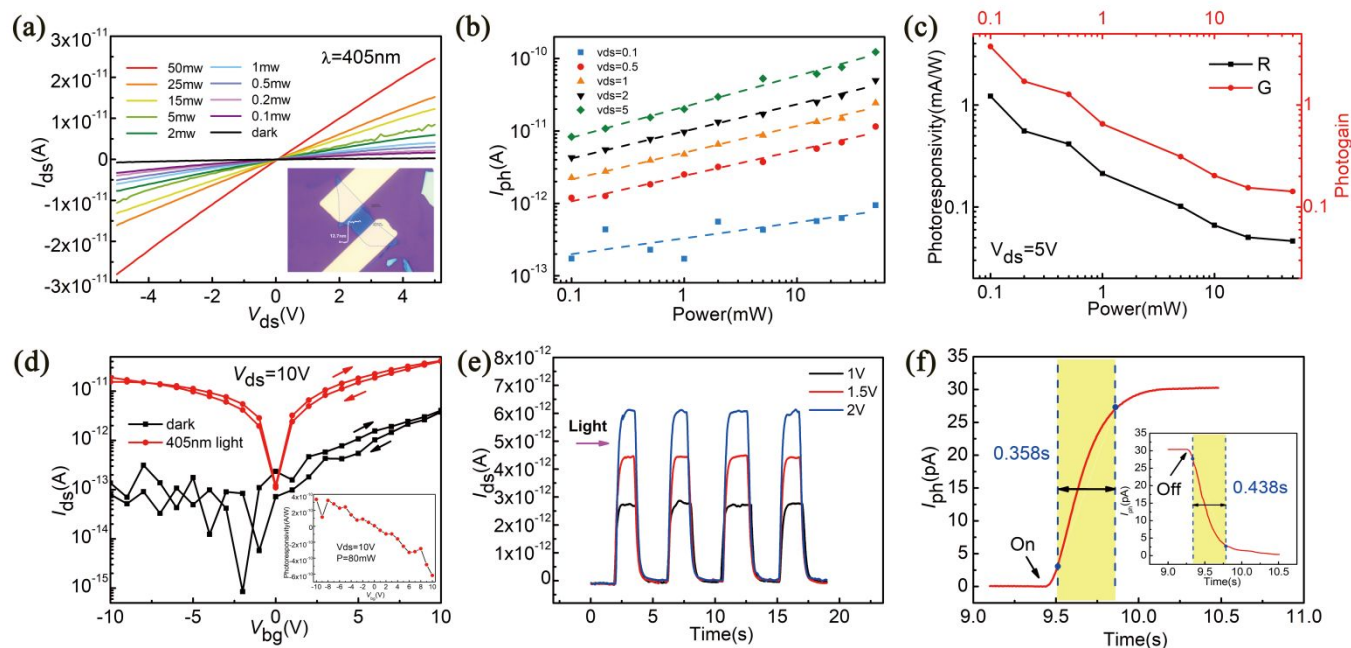


Fig 3 | Optical properties of the SnP<sub>2</sub>S<sub>6</sub> photodetector.

(a)  $I_{ds}$ - $V_{ds}$  curves of SnP<sub>2</sub>S<sub>6</sub> under different light powers. Inset: The optical image of the device with thickness of 12.7 nm. (b) Photocurrent ( $I_{ph}$ ) vary with light power and  $V_{ds}$ . (c) Photoresponsivity (R) and photogain (G) vary with light power.  $R_{max} = 1.22$  mA/W,  $G_{max} = 3.74$ . (d)  $I_{ds}$ - $V_{bg}$  curves at dark and under 405 nm light. The dark current could be reduced to pA. Under light conditions, current hysteresis box becomes little. Inset: Photoresponsivity vary with  $V_{bg}$ , under  $V_{ds} = 10$  V,  $P = 80$  mW. (e) The ON/OFF test of photocurrent under 532 nm light and the bias from 1 V to 2 V. The ON-OFF ratio  $> 10^3$ . (f) The exponential fitting of the dynamic response of photocurrent for the rise and decay time. The deduced rise ( $\tau_{rise}$ ) and decay ( $\tau_{decay}$ ) time constants are about 0.358 s and 0.438 s.

The gate voltage dependence of the photoresponse of SnP<sub>2</sub>S<sub>6</sub> is also investigated. The transfer curves under light on and off conditions are shown in Fig.3(d), where an incident illumination power of 80 mW is applied. The curves show an increase of drain current by near two orders of magnitude once the device is illuminated, indicating that the photocurrent dominates in the entire operating range of the device. As can be seen from Fig.3(d), the photocurrent shows an ambipolar behaviour with a large gate-tunability.

In order to study the stability and photoresponse speed of the SnP<sub>2</sub>S<sub>6</sub> photodetector, the time-current curves with the 532 nm laser switching on and off was measured. For three different bias voltages (1V, 1.5V and 2V), as shown in Fig.3(e), the time dependent photocurrent varies periodically with regular illumination, suggesting a stable and reversible photoresponse. The response speed is evaluated by a typical rise time/decay time, which is defined as the time over which the photocurrent increases from 10% to 90% (or decreases from 90 to 10% analogously) when the laser is switched on and off. Fig.3(f) shows the result of rise time and decay time of  $\tau_{rise} = 0.358$  s and  $\tau_{decay} = 0.438$  s for  $V_{ds} = 2$  V, where the rising and falling parts of the curve can be fitted with an exponential function  $I = I_0 + Ae^{t/\tau_{rise \text{ or } decay}}$ , where  $\tau_{rise \text{ or } decay}$  is rise or decay time constant.

### Dielectric properties of SnP<sub>2</sub>S<sub>6</sub> film

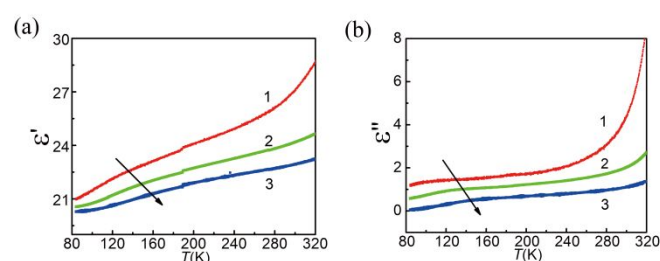


Fig. 4 | Temperature dependence of SnP<sub>2</sub>S<sub>6</sub> crystal dielectric permittivity real (a) and imaginary (b) parts at frequencies 1(1 kHz), 2(10 kHz), 3(40 kHz).

We first studied the temperature dependence of the dielectric permittivity at different frequencies. For this, samples of SnP<sub>2</sub>S<sub>6</sub> were used, obtained from the gas phase with a size of  $5 \times 5 \times 0.05$  mm<sup>3</sup>. Silver paste deposited on opposite planes of the test sample was used as electrodes. The measurements were carried out in the frequency range 10 Hz to 50 kHz and in the temperature range 80 to 400 K. For research, we used a GW INSTRON LCR-819 meter and an immersed to liquid nitrogen cryostat with a temperature measurement accuracy of 0.001 K. The cooling rate was 0.1 K/min at a measuring field strength of 1 V/cm. As can be seen in Fig.4(a), (b), when the SnP<sub>2</sub>S<sub>6</sub> crystals are cooled, a monotonic decrease in the permittivity (a) and dielectric losses (b) are observed. At high temperatures, low frequency  $\epsilon'$  and  $\epsilon''$  increase strongly, which is

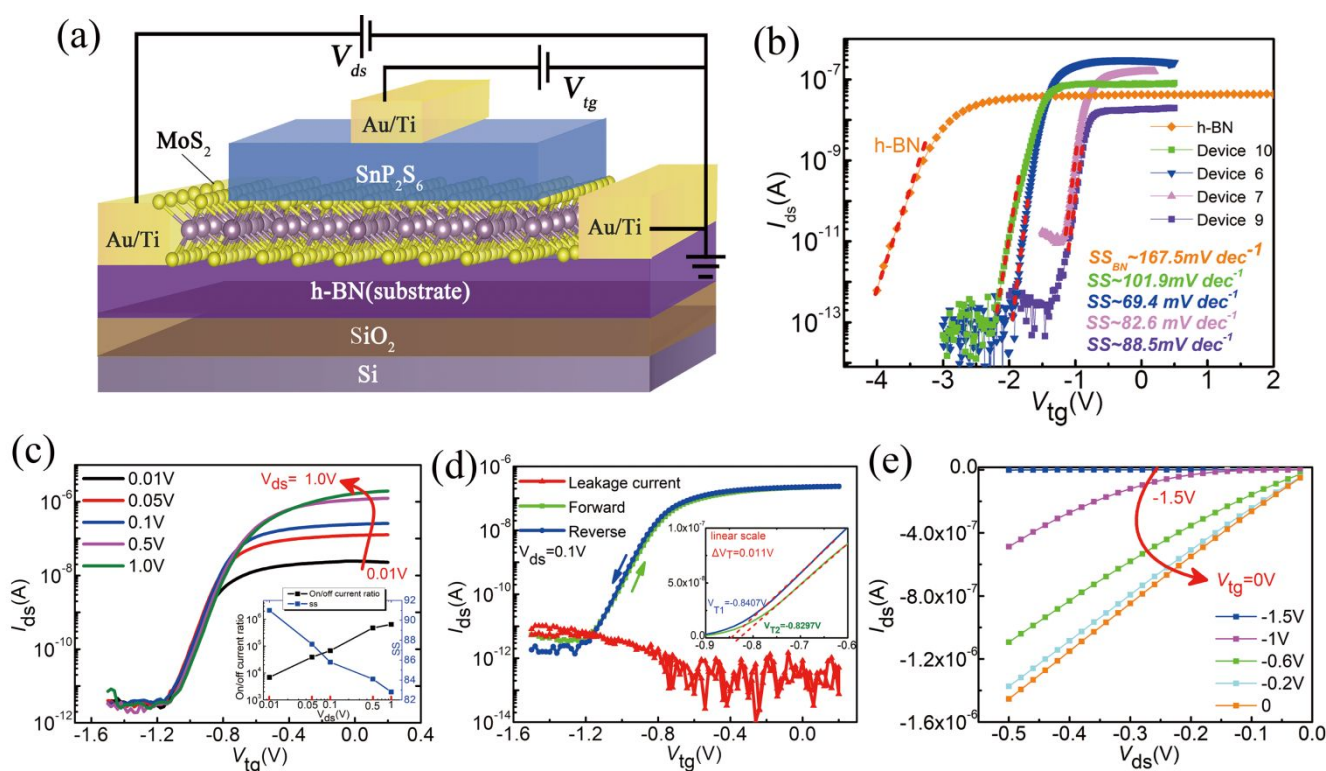
most likely due to an increase in the conductivity of the sample. The dielectric value of  $\text{SnP}_2\text{S}_6$  ( $\approx 23$ ) is much higher than that of h-BN ( $\approx 5.0$ ), which is comparable with or higher than those of traditional high-k oxide dielectrics such as  $\text{Al}_2\text{O}_3$ <sup>45</sup>,  $\text{HfO}_2$ <sup>46</sup>,  $\text{ZrO}_2$ <sup>47</sup>.

Considering that  $\text{SnP}_2\text{S}_6$  film has excellent insulating properties, good chemical stability and high dielectric constant, it shows great potential for integration into 2D FETs as a dielectric layer to modulate carrier density in a semiconducting channel. To demonstrate its advantages in this regard, we fabricate 2D FETs and evaluate the gating effect of  $\text{SnP}_2\text{S}_6$  as top-gate dielectrics.

The well-studied 2D semiconductor  $\text{MoS}_2$  was chosen as a representative channel material. Initially, we build  $\text{MoS}_2/\text{SnP}_2\text{S}_6$  FETs on  $\text{Si}/\text{SiO}_2$  (corresponding schematic is provided in **Supplementary fig.S5A**), but the result shows that the hysteresis gap of the devices is too big (**Supplementary fig.S5C**). Here, in order to reduce the hysteresis gap, we use h-BN as the substrate, which can be prepared with a very uniform thickness and an atomically flat surface<sup>48</sup>, and this method is proved to improve the performance of  $\text{MoS}_2$  FETs as device substrate<sup>49, 50</sup>. The schematic illustration of our device is shown in **Fig.5(a)** (Optical image in **Supplementary fig.S6** and **fig.S7D**).

The thin-film 2D materials are obtained through a mechanical exfoliation process (see **Supplementary fig.S7A, B, C**). Through the process described in the methods, 7 devices based on  $\text{SnP}_2\text{S}_6/\text{MoS}_2$  flakes of various thicknesses were fabricated. The optimal electric performance was observed in the devices with  $\text{SnP}_2\text{S}_6$  of about 49.7 nm thickness (details in **Supplementary fig.S6** and **fig.S8**).

Meanwhile, a top-gate h-BN/ $\text{MoS}_2$  FET was also fabricated for comparison according to the schematic in **Supplementary fig.S9C**, in which h-BN (with thickness of about 23.44 nm) is thinner than  $\text{SnP}_2\text{S}_6$  (Optical and AFM image in **Supplementary fig.S6**). **Fig.5(b)** shows the top-gate transfer ( $I_{ds}-V_{tg}$ ) characteristics measured for h-BN/ $\text{MoS}_2$  FET (the orange points) and four  $\text{SnP}_2\text{S}_6/\text{MoS}_2$  FETs on insulated h-BN substrate (see **Supplementary fig.S7D**). For the former one, a subthreshold slope (SS) value of  $167.5 \text{ mV dec}^{-1}$  and ON/OFF ratio of  $10^5$  was obtained. While for the latter,  $\text{SnP}_2\text{S}_6/\text{MoS}_2$  FETs, a small subthreshold slope (SS) value as low as  $69.4 \text{ mV dec}^{-1}$  for device 6 (Blue points) is achieved under a drain bias of  $0.1 \text{ V}$ , which is close to the limit value<sup>51</sup> of  $60 \text{ mV dec}^{-1}$ . The device could even reach high ON/OFF ratio of  $10^7$  and high carrier mobility of  $28.498 \text{ cm}^2\text{V}^{-1}\text{s}^{-1}$  at higher drain bias (**Supplementary fig.S10 (A)**). The threshold voltage ( $V_t$ ) is a parameter used to describe how much voltage is needed to initiate the channel conduction, which of the fabricated device is as low as  $-1.1 \text{ V}$  (**Supplementary fig.S10 (B)**). The FETs with smaller SS value, low on voltage and low threshold voltage can achieve effective switching at a low  $V_{ds}$ , thereby reducing the leakage and power consumption of FETs devices. It's remarkable for our  $\text{SnP}_2\text{S}_6$  to have such an excellent SS. In order to investigate the reason, we have carried out the following discussions. Through the contrast above, it is mostly possible that thicker  $\text{SnP}_2\text{S}_6$  films of about 47.7 nm to 49.7 nm show best performance with greater possibility of applications. What's more, they have stronger gate control capability and low leakage current than those of 23.44 nm-thick h-BN.



**Fig. 5 | Electronic properties of top-gated  $\text{MoS}_2$  FETs using  $\text{SnP}_2\text{S}_6$  as the gate dielectric. (a)** Schematic of  $\text{SnP}_2\text{S}_6/\text{MoS}_2$  FETs using h-BN as substrate. **(b)**  $I_{ds}$ - $V_{tg}$  curves of  $\text{MoS}_2$  FETs at 300 K. The orange points represent h-BN/ $\text{MoS}_2$  FET as contrasting device, in which h-BN is thinner than  $\text{SnP}_2\text{S}_6$ . Others are  $\text{SnP}_2\text{S}_6/\text{MoS}_2$  FETs which exhibit subthreshold slope (SS) down to **69.4 mV/dec**. **(c-e)** show the best performance of

SnP<sub>2</sub>S<sub>6</sub>/MoS<sub>2</sub> FET (Device 7). (c)  $I_{ds}$ - $V_{tg}$  characteristics measured at drain bias from 0.01 V to 1 V, the inset shows ON/OFF current ratio and SS extracted at different  $V_{ds}$ . (d) Double-sweep  $I_{ds}$ - $V_{tg}$  curves show small hysteresis and leakage current, the inset is the current in linear scale (hysteresis gap  $\Delta V_T \approx 0.011$  V). (e)  $I_{ds}$ - $V_{ds}$  characteristics measured at different  $V_{tg}$ .

which means larger capacitance value than h-BN in MoS<sub>2</sub> FETs ( $C_{SnP_2S_6} > C_{h-BN}$ ). And since SnP<sub>2</sub>S<sub>6</sub> is thicker than h-BN ( $d_{SnP_2S_6} > d_{h-BN}$ ), according to the equation

$$C = \epsilon_0 \epsilon_r / d,$$

where C is capacitance value per unit area of the gate dielectric,  $\epsilon_0$  is permittivity of vacuum and  $\epsilon_r$  is the relative dielectric constant. This confirms the high dielectric constant of SnP<sub>2</sub>S<sub>6</sub> and advantages of integrating SnP<sub>2</sub>S<sub>6</sub> nanosheet with semiconductor crystals to design FETs compared with h-BN, demonstrating the potential of 2D SnP<sub>2</sub>S<sub>6</sub> as high-k dielectric for low power and miniaturized electronics.

As shown in Fig. 5(c,d), the typical SnP<sub>2</sub>S<sub>6</sub>/MoS<sub>2</sub> device exhibits high ON/OFF ratio and steep SS simultaneously. The effect of the increasing drain bias, resulting in augmented drive current and steeper SS value, are displayed in Fig. 5(c). It can be seen that the best transistor performance is achieved at  $V_{ds} = 1$  V, with a maximum measured on-state current of  $\sim 1 \mu A$ , ON/OFF current ratio close to  $10^6$ , SS value as small as  $82.6 \text{ mV dec}^{-1}$  (inset of Fig. 5(c)) and the leakage current as low as  $\sim 10^{-11} \text{ A}$  (Fig. 5(d)), respectively. The

corresponding hysteresis stability of the device was demonstrated in Fig. 5(d). To describe the hysteresis stability, we introduce the hysteresis gap ( $\Delta V_T$ ), which is defined as the difference between the forward and reverse sweep threshold voltage<sup>52</sup>. As obtained in the linear scale in the inset of Fig. 5(d),  $\Delta V_T = |V_{T1} - V_{T2}| \approx 0.011 \text{ V}$ , was determined as low as  $0.011 \text{ V}$ , which is a fairly small value in FETs. However, the hysteresis gap is larger after two days in the glove box (Supplementary fig. S11D, E, F) and different hysteresis behaviours were observed in some of our devices (Supplementary fig. S11A, B, C). The hysteresis variation may be caused by flake quality or fabrication process in laboratory environment, which leads to charge trapping at the interface, adsorption/desorption of ambient molecules or atmospheric moisture in the exposed active layers<sup>50,53</sup>. The small hysteresis could be indicated that the fabricated SnP<sub>2</sub>S<sub>6</sub>/MoS<sub>2</sub> device shows atomic flat and sharp interface. The corresponding output ( $I_{ds}$ - $V_{ds}$ ) characteristics measured for different  $V_{tg}$  (Fig. 5(e)) show a linear  $I_{ds}$ - $V_{ds}$  characteristics at low drain bias and saturated current at high drain bias.

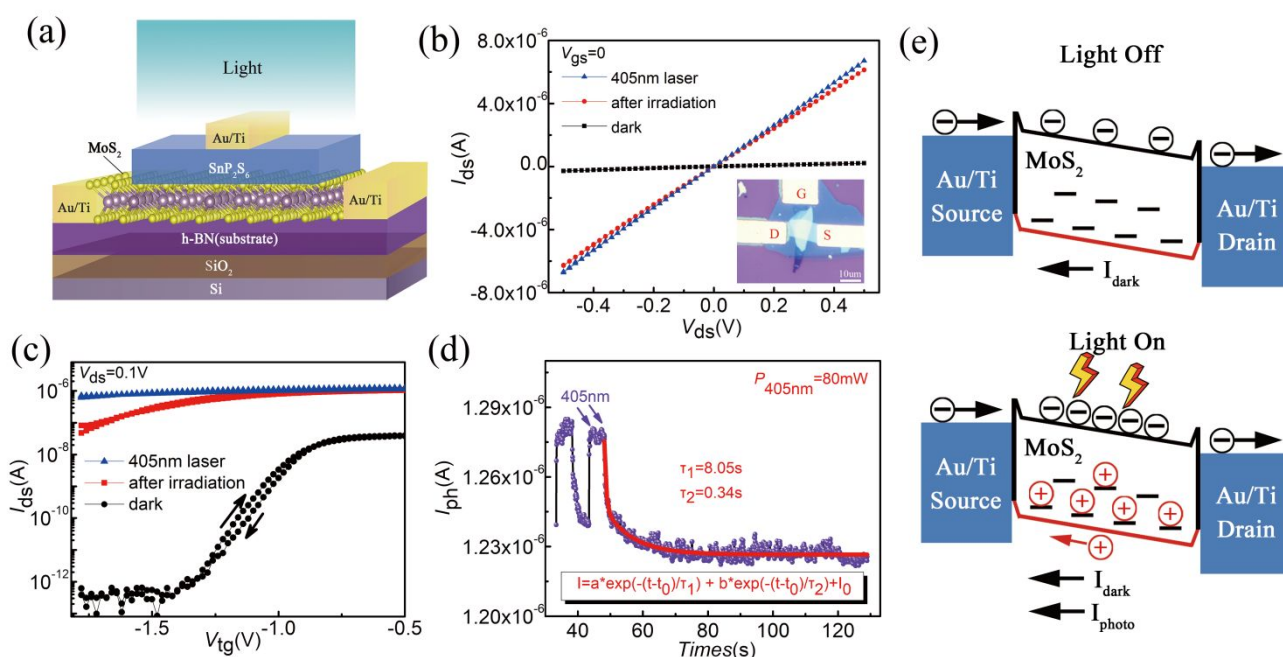


Fig. 6 | **Optical properties of top-gated SnP<sub>2</sub>S<sub>6</sub>/MoS<sub>2</sub> FETs.** (a) Schematic diagram of photogating dominated SnP<sub>2</sub>S<sub>6</sub>/MoS<sub>2</sub> FETs (Device 9). (b) Typical  $I_{ds}$ - $V_{ds}$  characteristics of Device 9 at dark (Black), under illumination (Blue) and after irradiation (Red), the inset is optical image of the FET. (c)  $I_{ds}$ - $V_{tg}$  characteristics shows transfer curve shifts in negative  $V_{tg}$  direction under and after irradiation. (d) Photocurrent decay after a longer light pulse with longer ( $\tau_1 \sim 8.05 \text{ s}$ ) and shorter ( $\tau_2 \sim 0.34 \text{ s}$ ) time constants, which corresponds to (b) and (c). (e) Band diagram of the device without (top panel) and with light (bottom panel), showing hole trapping in shallower and deeper trap states close to the valence band.

Breakdown electric field ( $V_{bd}$ ) is another important parameter to evaluate the strength of insulating layer. To test the current leakage and breakdown threshold, SnP<sub>2</sub>S<sub>6</sub> is sandwiched between two electrodes to perform a vertical device as shown in the inset of

fig.S12. At room temperature,  $V_{bg}$  are fixed at 0 V, and  $V_{ds}$  is swept<sup>54</sup>. The  $I_{ds}$ - $V_{ds}$  curve is plotted in fig.S12, where the initially low current increases slowly as the  $V_{ds}$  increase. The current increases dramatically as  $V_{ds}$  approaches  $V_{bd}$ . After breakdown, the device is

permanently destroyed as a large current flows across  $\text{SnP}_2\text{S}_6$ .  $V_{\text{bd}}$  is estimated to be about 5 V (with thickness of 26.5 nm), and the corresponding breakdown electric field  $F_{\text{bd}}$  could be reach as high as  $1.9 \text{ MV cm}^{-1}$ .

Furthermore, to investigate the optoelectronic properties of our FETs, we put one of our  $\text{SnP}_2\text{S}_6/\text{MoS}_2$  FETs on h-BN substrate (**Device 9**, its optical image in the inset of **Fig.6(b)**) under 405 nm laser irradiation of 80 mW. The output and transfer characteristics at dark (Black), under 405 nm irradiation (Blue) and after irradiation (Red) are separately depicted in **Fig.6(b)** and **Fig.6(c)** separately. Under illumination, the output characteristic of the device remains linear while reducing the electrical resistance of the semiconductor caused by the photoconductive effect (PC effect), which is originated from the extra free carriers generated by the photon absorption (**Fig.6(b)**). Meanwhile, as shown in **Fig.6(c)**, significant change in threshold voltage was observed when luminated by 405 nm light. The shifting of threshold voltage to the negative  $V_{\text{tg}}$  side is known as photogating effect (PG effect). PG effect is a phenomenon that photogenerated holes get trapped and then attract more electrons, thus gradually increasing the channel conductivity<sup>55, 56</sup>.

In **Fig.6(c)**, compared with the initial device in dark condition, the Off-state current after irradiation is dramatically increased by more than five-orders of magnitude at the same  $V_{\text{tg}}$ , implying a long-time decay of photoconductive state. As shown in **fig.S13**,  $\text{SnP}_2\text{S}_6$  has little effect on the optical response and persistent photocurrent of the device. Therefore, to explore the photocurrent decay of the device, we then check the transient characteristics by measuring the source-drain current at fixed  $V_{\text{tg}} = 0 \text{ V}$  in **Fig.6(d)**. An abrupt increase in  $I_{\text{ds}}$  is observed as soon as illuminated under 405 nm laser, as a result of the photoexcitation, followed by a much slower current increase, namely the PG effect. A very slow decay of the current can be observed when the light was switched off. This is because the trapped charges in PG effect cause long sustained conductivity<sup>56</sup>. In order to quantify the trapping time constant, transient data was fitted with a biexponential equation of the type<sup>57</sup>

$$I = I_0 + Ae^{-(t-t_0)/\tau_1} + Be^{-(t-t_0)/\tau_2}$$

where  $\tau_1$  and  $\tau_2$  are two relaxation time constants, possibly corresponding to two different photoresponse mechanisms. The fast part is due to photoconductivity effect. The slow part could be regarded with the defect trapping or photogating effect.  $I_{\text{ds}}$  decays exponentially from  $\sim 1.27 \mu\text{A}$  in time and saturated after 35 s with longer time constant  $\tau_1 \sim 8.05 \text{ s}$  and shorter time constant  $\tau_2 \sim 0.34 \text{ s}$ . (See the detailed parameter values in **Supplementary Table 2** and time constants for different wavelengths in **Supplementary fig.S14**). The shorter decay component could be attributed to PC effect and shallower traps, while the longer decay component might be due to PG effect caused by deeper (middle gap) traps in **Fig.6(e)**, as elucidated by Furchi et al<sup>55</sup>. Simultaneously, the longer time constants of the devices on Si/SiO<sub>2</sub> FETs are  $\sim 6.814 \text{ s}$  and  $\sim 13.39 \text{ s}$ , which are closed to 8.05s in the FETs on h-BN (See Details in **Supplementary fig.S15**). Consequently, we infer that the long persistent photocurrent measured in our devices is not likely due to the traps at the  $\text{MoS}_2/\text{SiO}_2$  or  $\text{MoS}_2/\text{h-BN}$  interface since our

$\text{SnP}_2\text{S}_6/\text{MoS}_2$  FETs fabricated on h-BN which have very less charge trapping<sup>50</sup>. So it is conjectured that the defect trapping might be mainly originated from photo charge trapping associated with  $\text{MoS}_2$  defects and the  $\text{MoS}_2/\text{SnP}_2\text{S}_6$  interface, which causes the PG effect. Therefore, improving the quality of  $\text{MoS}_2$  and optimizing the  $\text{MoS}_2/\text{SnP}_2\text{S}_6$  interface can be an approach to decrease this phenomenon.

Finally, the good stability, reliability and reproducibility of the devices have been summarized as shown in **fig.S16, 17,18, 19**.

To compare the performance of  $\text{SnP}_2\text{S}_6$  dielectric, the main set of achievements of the present study, such as dielectric constant, SS value and ON/OFF ratio were collected in Table 1. Compared with other new 2D dielectric materials,  $\text{SnP}_2\text{S}_6$ 's high performance is comparable to or even better than those of these materials.

**Table 1.** Comparison of performance based on  $\text{SnP}_2\text{S}_6$  with other new 2D dielectrics.

Materials	Dielectric constant	SS value [mV/dec]	ON/OFF ratio	Reference
Ta <sub>2</sub> O <sub>5</sub>	15.5	61	10 <sup>6</sup>	[58]
Bi <sub>2</sub> SeO <sub>5</sub>	21	75	10 <sup>5</sup>	[59]
CaF <sub>2</sub>	8.43	90	10 <sup>7</sup>	[53]
In <sub>2</sub> P <sub>3</sub> S <sub>9</sub>	24	88	10 <sup>5</sup>	[9]
Sb <sub>2</sub> O <sub>3</sub>	11.5	64	10 <sup>8</sup>	[60]
$\text{SnP}_2\text{S}_6$	23	69.4	10 <sup>7</sup>	This work

## Conclusion

In this paper, we have systematically investigated the photoelectric and electrical properties of the novel 2D  $\text{SnP}_2\text{S}_6$  crystal. Moreover, it was integrated into  $\text{MoS}_2$  FETs to serve as top gate dielectric. The results demonstrate its excellent dielectric properties including large dielectric constant ( $\approx 23$ , compared with h-BN), a fairly small SS value (as low as  $69.4 \text{ mV dec}^{-1}$ ), high ON/OFF ratio (reaching 10<sup>7</sup>), negligible hysteresis, relatively high breakdown voltage ( $\approx 1.9 \text{ MV cm}^{-1}$ ). Due to its semiconductor nature, good photoelectric properties and good reliability and reproducibility (0.5 standard deviation of threshold voltage),  $\text{SnP}_2\text{S}_6$  could become a potential candidate to construct advanced optoelectronic and electronic devices.

## Methods

**Material Characterization.** Raman spectra is obtained in a Horiba confocal Raman system with an excitation laser of 532 nm. The morphology characterization and thickness calibration of the  $\text{SnP}_2\text{S}_6$  film as well as thickness measurement of h-BN and  $\text{MoS}_2$  are carried out with optical microscope and AFM instruments.

**Fabrication of  $\text{SnP}_2\text{S}_6$  phototransistor and  $\text{SnP}_2\text{S}_6/\text{MoS}_2$  FET devices.** The  $\text{SnP}_2\text{S}_6$  phototransistor was fabricated by exfoliating the thin flakes onto heavily doped silicon substrates. The top electrodes are defined using standard photolithography process followed by



thermal evaporation of the Ti/Au metal, and lift-off process. We have built three batches of SnP<sub>2</sub>S<sub>6</sub>/MoS<sub>2</sub> devices on Si/SiO<sub>2</sub> substrate and on h-BN substrate, respectively. Thin-film 2D materials are obtained through a mechanical exfoliation process first. Then the 2D nanosheets of top-gate FETs are stacked up by polydimethylsiloxane (PDMS) assisted transfer technique onto h-BN nanosheet to develop heterostructures. And the source-drain and gate electrodes Ti/Au of every device were patterned and deposited by photolithography and a lift-off process. Preparation methods for other devices including back-gate FETs and top-gate FETs on different substrates vary according to their schematic.

**Electrical and optical characterization measurements.** We tested the device directly after fabrication. Electrical and optical characterization of devices consisted of I<sub>D</sub>-V<sub>G</sub>, I<sub>D</sub>-V<sub>D</sub> and I<sub>ph</sub>-t characteristics. All measurements are performed under vacuum conditions at room temperature (300 K) using probe station.

### Author contributions

J. H. and A. Z. carried out the electrical and optoelectronic measurement. E. P., J. C. and R. B. assistant the device fabrication and electrical measurement. E. P. carried out the AFM measurement. J. L. and X. J. carried out the Raman characterization. A. M. and Y. V. conducted dielectric measurement and crystal growth. Q. L., G. C. and P. M. analyzed the data. J. H., A. Z. and F. L. wrote the manuscript. Q. L. and G. C. and revised the manuscript. F. L. supervised the project. All the authors comments on the manuscript.

### Conflicts of interest

There are no conflicts to declare.

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